

# The characteristics of nickel-based elliptical cylinder plasmonic nano-waveguides in the near-infrared wavelength

Scientific research paper

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### 1 Introduction

 The extraordinary advantages of surface plasmons (SPs) in compressing light into regions much smaller than the diffraction limit, have made them conceivable candidates for the manufacture and confinement of optical coordinate circuit components [1-5]. In some previous waveguides, noble metals [6] have been properly utilized for SPs [7] modes, including nanowires [8-13], wedge waveguides [14], waveguides with dielectric charge [15], and etc. Nevertheless, they are still involved with some problems such as field limits. To solve this problem, the hybrid plasmonic waveguides (HPWs) [16-32] have been utilized instead of waveguides based on noble metals [33-35], with the hope that they would result in less dissipation and more field confinement. However, in practice, the desired results have not been obtained, so it is still of vital importance for researchers to achieve desired results in plasmonics. In several studies that we will refer to in the following, the characteristics of the sodium metal have been examined for their applications in plasmonic waveguides. In a plasmonic nanolaser based on sodium metal, Yang and colleagues [36] found that sodium has incorporated a lower diffraction limit at near-infrared wavelengths. In another study, Tao and his colleagues [37,38] pinpointed the existence of directional sets in plasmonic waveguides based on sodium metal, and Da Teng and his colleagues [39] proposed two sodium nanowires based on sodium metal with two layers of dielectric coating. Comparing them with silver metal, Da Teng and his colleagues observed that the nanowire Na shows less dissipation and more grounded field confinement. Recently a palladium-based elliptical cylinder plasmonic waveguide (PECPW) with two dielectric layers as a coating in different states has been taken under consideration [40]. They showed that the highest value for the propagation length  $L_{\rm{snn}}$  and the

figure of merit FOM of this waveguide is for a case of single palladium nanowire with  $b=4a$  and without any dielectrics. In this paper, we propose a plasmonic nanowaveguide based on nickel metal within the near infrared range (NIR) run in which longer propagation distances exist, where better figure of merit (FOM) and stronger confinement are expected in comparison to the similar plasmonic waveguides (e.g., see [40, 41]). In this paper, the characteristics of the suggested nanowaveguide have been examined based on the finite element method and in accordance with the wavelength and various thicknesses of the dielectric layers. The results demonstrate that a cylinder elliptical plasmonic waveguide based on nickel metal has a longer propagation length, lower dissipation and smaller ordinary mode region compared to sodium metal. However, the results are not as promising as those of PECPW. As a result, nickel metal can be utilized in nanophotonic circuits.

## 2 Waveguide design and mathematical formulation



Figure 1. The cross-section of the proposed NECPNW.

Figure 1, shows the cross-section of the plasmonic waveguide structure based on nickel metal. In this structure an elliptical cylinder nanowire with two dielectric layers ( $MgF_2$  and Ge) is used. As you can see,  $a$  and  $b$  are respectively, the semi-minor axis and semimajor axis of the proposed waveguide cross section.  $t_{\text{MgF}_2}$  is the thickness of the MgF<sub>2</sub> dielectric layer with a relative permittivity of  $\varepsilon_{\text{MgF}_2} = 1.7$  [39] and  $t_{\text{Ge}}$  is the thickness of the Ge dielectric layer with a relative permittivity of  $\varepsilon_{\text{Ge}} = 16.2$  [39]. Investigation and analysis of surface plasmon polaritons propagation mode in NECPNW structure is done in very small dimensions for wavelength 400 − 1000 nm. In addition, it is possible to calculate the dielectric function

of nickel metal using the Drude-Lorentz modal, which is expressed as [41]:

$$
\varepsilon_{DL}(\omega) = \varepsilon_1 + i\varepsilon_2
$$
  
=  $\varepsilon_b - \frac{\omega_p^2}{\omega^2 + i\omega\gamma_p}$   
+  $\frac{f_1\omega_1^2}{\omega_1^2 - \omega^2 - i\omega\gamma_1}$ . (1)

The parameters used in calculating the dielectric function of nickel are according to Table 1, [42]:

Table 1. Drude-Lorentz model coefficients used for nickel [42].

Parameter	Value
f <sub>1</sub>	0.137
$\omega_1$	0.415 eV
$\gamma_1$	$0.765 \text{ eV}$
$\boldsymbol{\omega_p}$	13.28 eV
$\gamma_p$	0.038 eV

 The effective refractive index of the proposed waveguide is calculated as [40]:

$$
n_{\text{eff}} = \frac{k_z}{k_0},\tag{2}
$$

where  $k_0$  is the wave number in vacuum and  $k_z$  is the component of z the wave number. The complex effective mode index is solved as [7]:

$$
n_{\text{eff}} = n_{\text{eff}-\text{r}} + i n_{\text{eff}-\text{i}} \tag{3}
$$

In addition, the propagation length is obtained from equation  $(4)$ ,  $[40]$ 

$$
L_{spp} = \frac{1}{k_0 \text{Im}(n_{\text{eff}})},
$$
\n(4)

 $A_n$ ,  $A_{\text{eff}}$  and  $A_0$  respectively, represent the normalized mode area, the effective mode area, and the diffraction limited mode region in the vacuum environment, and their relationship is as [40]:

$$
A_{n} = \frac{A_{\text{eff}}}{A_{0}} = \frac{1}{A_{0}} \frac{\int_{-\infty}^{+\infty} W(x, y) dx dy}{\max[W(x, y)]},
$$
 (5)

So that,  $W(x,y)$  is the energy density of waveguide plasmon mode [40] and is defined as:

$$
W(x, y)
$$
  
=  $\frac{1}{2}Re\left\{\frac{d[\varepsilon_0 \varepsilon_r(x, y)\omega]}{d\omega}\right\} |E(x, y)|^2$   
+  $\frac{1}{2}\mu_0 |H(x, y)|^2$ . (6)

Finally, the waveguide figure of merit criterion is [43] calculated by:

$$
FOM = \frac{L_{\rm SPP}}{\sqrt{\frac{A_{\rm eff}}{\pi}}}.
$$
\n(7)

 The Results of this research have been obtained using the finite element method in the COMSOL-Multiphysics 6.0 software based on the mentioned parameters and variables.

3 Results and discussionFigure 2, shows how the electric field is confined in the NECPNW structure when no dielectric is used where  $a = 50$  nm and different values of  $b=a$ , 2a, 4a are predicted for the main radius and wavelength of  $\lambda = 400 - 1000$  nm. As can be seen in Fig.  $2(a)$ , in the case where b=a, a relatively strong electric field is created around the surface of the nickel metal, and with increasing the value of b, that is, for Figs.  $2(b)$  and  $2(c)$  in which the semi-major axis of  $b = 100$  nm and  $b = 200$  nm are considered respectively, the electric field is created on both sides of  $b$  and this quantity increases with the increase of  $b$ . Figure 2(d) shows that the diagram of the electric field around the nanowaveguide is drawn in the direction of the y axis according to the mentioned structure.

 Figure 3 illustrates the results of the propagation of the plasmonic nanowaveguide based on nickel metal in the case where no dielectric is used. In this structure  $t_{\text{MgF}_2} = t_{\text{Ge}} = 0$  and the semi-minor axis of the nanowire section  $a = 50$  nm and three different values the semi-major axis of the nanowire section  $b =$ 50,100, 200 nm are considered. In Fig. 3(a) it can be seen that the real part of effective mode index decreases with the increase of  $\lambda$  and b. Figure 3(b) shows that the propagation length has increased with the increase of wavelength and the semi-major axis, and this nanowaveguide has resulted in a longer propagation length compared to the waveguide based on sodium metal. Figure 3(c) shows that the normalized mode area for a given quantity of  $A_n$  of the nanowire for the semimajor decreases with the increase of the wavelength and for different values of the semi-major axis at a given wavelength it also increases with the increase of b. Figure 3(d) shows the changes of the nanowire figure of merit in terms of wavelength, which increases with the increase of the wavelength and the size of the semimajor axis.



Figure 2. Distribution of the electric field in the cross-section of the elliptical cylinder nanowire at (a)  $b=a$ , (b)  $b=2a$ , (c)  $b=4a$ , and (d) the electric field distribution in y-direction.

 Figure 4 shows the distribution of the electric field around the nickel nanowire for the elliptical cylinder nanowave with MgF<sub>2</sub> dielectric coating for  $a = 50$  nm and  $b = 4a$  and different values of  $t_{\text{MgF}_2}$ . As can be seen in Figs. 4(a), 4(b), and 4(c), the thickness of the dielectric coating is considered to be 10, 15, and 20 nm, respectively, and the intensity of the modal field on both



Figure 3. Modal properties nanowire in vacuum without dielectric coating for different values of semi-major axis. (a)  $Re(n_{eff})$ , (b)  $L_{\rm{snn}}$ , (c)  $A_n$  and (d) FOM.

sides of the semi-major axis of the nanowire decreases with the increase of the dielectric thickness. In Fig. 4(d), the diagram of the modal field around the nanowaveguide for the above values is drawn in the direction of the y axis.



Figure 4. Distribution of the electric field in the cross-section of NECPNW with dielectric coating MgF<sub>2</sub> and  $b=4a$ . (a)  $t_{\text{MgF}_2}$  = 10 nm, (b)  $t_{\text{MgF}_2} = 15$  nm, (c)  $t_{\text{MgF}_2} = 20$  nm and (d) distribution of the electric field in the cross-section of NECPNW for different values of the  $t_{\text{Mg}_{2}}$ .

 The different parts of Fig. 5 show the behavior of  $Re(n_{eff})$ ,  $L_{spp}$ ,  $A_n$  and FOM in terms of wavelength, which are plotted for  $a = 50$  nm,  $b = 4a$ ,  $t_{Ge} = 0$  and different values of  $t_{\text{Mg}_2}$ . So that  $t_{\text{Mg}_2}$  = 10, 15, 20 nm. As can be seen in Fig. 5(a), the value of  $Re(n_{eff})$  for a certain  $t_{MgF_2}$  decreases with the increase of the wavelength. In addition, this value for a certain a wavelength decreases with the increase of the dielectric thickness. Figure 5(b) shows that the propagation length  $L_{\rm spp}$  increases with the increase of the wavelength and decreases with the increase of the thickness of  $MgF_2$  dielectric layer. It can be seen in Fig. 5(c) that  $A_n$  for a specific dielectric thickness decreases slightly with the increase of the wavelength and is of the order of  $10^{-2}$ , but for a certain wavelength, the process of  $A_n$ changes is different from the dielectric layer thickness.



Figure 5. Properties of nano-waveguide with constant consideration of semi-major axis  $b=4a$  and different values of  $t_{\text{MgF}_2}$ . (a)  $Re(n_{eff})$ ,(b)  $L_{spp}$ ,(c)  $A_n$  and (d) FOM.

 Finally, in Fig. 5(d) the graph of the dependence of the figure of merit FOM on the wavelength is drawn, which is similar to Fig.  $5(b)$ , the value of this quantity increases with the increase of the wavelength and decreases with the increase of the thickness of the dielectric layer.

Figure 6 shows the dependence of Re( $n_{\text{eff}}$ ),  $L_{\text{spp}}$ ,  $A_{\text{n}}$ and FOM on the thickness of the second dielectric layer of germanium with size  $t_{\text{Ge}} = 5 \text{ nm}$ , for different values of the semi-major axis of the elliptical cylinder nickel nanowire with sizes  $b = 50, 100, 200$  nm.



Figure 6. Dependence of nano-wave properties on  $t_{\text{Ge}} = 5 \text{ nm}.$  (a)  $Re(n_{eff}),$ (b)  $L_{\text{spp}},$ (c)  $A_n$  and (d) FOM.

It can be seen in Fig. 6(a) that the size of  $Re(n_{eff})$ increases with the increase of the thickness of the dielectric layer  $t_{\text{Ge}}$  and for a certain value of the thickness of the dielectric layer  $t_{\text{Ge}}$ , the size of Re( $n_{\text{eff}}$ )

decreases with the increase of the semi-major axis. In Figs. 6(b) and 6(d), it can be seen that the propagation length  $L_{\text{spp}}$  and the figure of merit FOM increase with the increase of the thickness of the dielectric layer  $t_{\text{Ge}}$ , and for a certain value of the thickness of the dielectric layer  $t_{\text{Ge}}$ , its size decreases with the increase of the semi-major axis of the nanowire cross-section. Finally, in Fig. 6(c), the dependence of  $A_n$  on the thickness of the dielectric layer  $t_{\text{Ge}}$  can be seen, its size decreases with the increase in the thickness of the dielectric layer, and for a value of the thickness of the dielectric layer, the value of  $A_n$  decreases with the increase of the semimajor axis of the nanowire cross-section. Hence, the proposed waveguide shows longer propagation length and better figure of merit than the similar plasmonic waveguides (e.g., see [40, 41]).

### 3 Conclusions

 In this article, a cylinder elliptical plasmonic nanowaveguide based on nickel metal was investigated in NIR range. The results of the investigation of this nanowaveguide show that in the case where  $b=4a$ ,  $t=$  $th = 0$  nm, in the range of  $\lambda = 400 - 1000$  nm, the propagation length  $L_{SDD} = 99.92726$  nm and the figure of merit  $FOM = 84.2883$  and the lowest value for the normalized mode area  $A_n = 0.00912$  is obtained. In the condition that  $b=4a$  and the thickness of the first dielectric layer, we considered the values of  $t_{\text{MgF}_2}$  = 10, 15, 20 nm, when  $t_{\text{MgF}_2} = 10$  nm, the propagation length  $L_{\rm spp} = 21.81906$  nm and  $FOM = 75.1905$  and  $A_n = 0.017$ , which had significant and better results than the other two modes. Finally, when we considered two layers of dielectric coating with thicknesses of  $t_{\text{Ge}} = 400 \text{ nm}$  and  $t_{\text{MgF}_2} = 5 \text{ nm}$  for the plasmonic nanowaveguide, the best case and results were obtained for  $b=a$ , so that  $L_{\text{spp}}=63.481668$  nm and  $FOM =$ 22.37486 and  $A_n = 0.002$ . The obtained results indicate that the ellipsoidal cylinder plasmonic waveguide based on nickel metal has not only lower losses, but also a measure of competence and longer propagation length and a stronger field enclosure than sodium. It shows that it can be used in highperformance infrared plasmonic devices.

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